

Stat Phys IV

Lecture 2

Reference "The Mathematics of Brownian Motion and Johnson noise"
Gillespie Am. J. Phys. 1996

(Continuous Markov Process Theory)

Mathematical Foundations of Langevin Eqs:

Ordinary calculus $\frac{dx(t)}{dt} = A(x(t), t)$ A : smooth fct.

can be expressed as: (where $x(t)$ is called PROCESS)

$$\boxed{x(t+dt) = x(t) + A(x(t), t) dt} \quad (1) \quad \text{"UPDATE FORMULA"}$$

Properties of process:

(1) DETERMINISTIC since $x(t+dt)$ is given UNEQUIVOCALLY

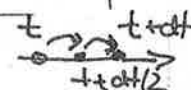
(2) CONTINUOUS since $x(t+dt) = x(t)$ for $dt \rightarrow 0$

(3) MEMORYLESS since $x(t+dt)$ depends on $x(t), t$.

$\Rightarrow$ CONTINUOUS, MEMORYLESS, DETERMINISTIC PROCESS.

Generalization to a CONTINUOUS, MEMORYLESS, STOCHASTIC PROCESS $\rightarrow$ CONTINUOUS MARKOV PROCESS.

Naïve approach: $x(t+dt) = x(t) + \sqrt{dt} A(x(t), t)$ (2)

This update formula is not consistent since, 

$$(1) \quad x(t+dt/2) = x(t) + A(x(t), t) \cdot \sqrt{dt/2}$$

$$(2) \quad x(t+dt) = x(t+dt/2) + A(x(t+dt/2)) \cdot \sqrt{dt/2}$$

inserting yields:

$$x(t+dt) = x(t) + A(x(t)) \sqrt{dt/2} + A(x(t+dt/2)) \sqrt{dt/2}$$

in the limit of $dt \rightarrow 0$

$$x(t+dt) = x(t) + 2 \sqrt{\frac{dt}{2}} A(x(t), t) \rightarrow \text{different from Equation (2)}.$$

The self consistency condition limits all continuous, memoryless stochastic processes to evolve under generalizations of (1) "LANGEVIN EQUATIONS".

Note. the condition of a memoryless process is not as restrictive as one may think: example:

e.g. reduce $m\ddot{x} = -\frac{\partial V}{\partial x}$ to 1st order diff eqs. $\begin{cases} m\dot{v} = -\frac{\partial V}{\partial x} \\ \dot{x} = v \end{cases}$ BIVARIATE
CONT. MEMOR.
PROCESS

Review (brief) Random Variables

$P(y)$ Probability density fct. of variable y ($P(y)dy$ is prob. to find $y \dots y+dy$.)

$$\langle f(y) \rangle = \lim_{N \rightarrow \infty} \frac{1}{N} \sum_{i=1}^N f(y_i) = \int_{-\infty}^{\infty} P(y) f(y) dy$$

VARIANCE: $\text{VAR}\{Y\} = \langle (Y - \langle Y \rangle)^2 \rangle = \langle Y^2 \rangle - \langle Y \rangle^2 \equiv \langle \Delta Y^2 \rangle$

STANDARD DEVIATION $\sigma \equiv \sqrt{\text{VAR}\{Y\}}$

COVARIANCE $\text{COV}\{Y_1, Y_2\} = \langle (Y_1 - \bar{Y}_1)(Y_2 - \bar{Y}_2) \rangle$
 $= \langle Y_1 Y_2 \rangle - \langle Y_1 \rangle \langle Y_2 \rangle$

for statistically independent variables: $\text{COV}\{Y_1, Y_2\} = 0$

GAUSSIAN RANDOM VARIABLE

$$P(y) = \frac{1}{\sqrt{2\pi}\sigma} \exp\left(-\frac{(y-\mu)^2}{2\sigma^2}\right) \quad Y = N(\mu, \sigma^2)$$

$\uparrow$ mean $\uparrow$ variance

Unit random variable: $N(0,1) \equiv N$

Properties: $\alpha + \beta N(\mu, \sigma^2) = N(\alpha + \beta\mu; \beta^2\sigma^2)$

thus for unit random variable:

$$\alpha + \beta N(0,1) = N(\alpha, \beta^2)$$

• $N(\mu_1, \sigma_1^2) + N(\mu_2, \sigma_2^2) = N(\mu_1 + \mu_2, \sigma_1^2 + \sigma_2^2)$

CENTRAL LIMIT THEOREM

$$\tilde{X} = \sum_{i=1}^k X_i^{(1)} + X_2^{(1)} + \dots + X_k^{(1)}$$

X_k : statistically independent but not normally distributed

as $k \rightarrow \infty$ leads to a NORMAL i.e. GAUSSIAN DISTRIBUTION.

DEFINITION: A STOCHASTIC PROCESS $X(t)$ is a random variable whose probability density depends on time t , $X(t_1)$ and $X(t_2)$ are two different random variables.

$X(t)$ is a CONTINUOUS, MEMORYLESS, STOCHASTIC PROCESS (i.e. A CONTINUOUS MARKOV PROCESS) if:

(i) MEMORYLESS i.e. $E(dt, x, t) = X(t+dt) - X(t)$

E : increment depends only on $t, dt, X(t)$

(ii) $E(x, t, dt)$ depends "smoothly" on t, x, dt

(iii) CONTINUOUS i.e. $E(x, t, dt) \rightarrow 0$ for $dt \rightarrow 0$.

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Conditions (i)-(iii) imply that a MEMORYLESS, STOCHASTIC, CONTINUOUS MARKOV PROCESS must have the form,

$$\boxed{E(t, x, dt) = \underbrace{A(x(t), t) dt}_{\text{Drift-term}} + \underbrace{D^{1/2}(x(t), t) N(t) \sqrt{dt}}_{\text{Diffusion}} \quad N(0,1)} \quad (3)$$

Proof requires as before to consider partial intervals ie $dt = \left(\sum_{i=1}^N dt^i \right)$ N : intervals of dt increment.

$$\begin{aligned} \langle N^2(t) \rangle &= 1 \\ \langle N(t) \rangle &= 0 \end{aligned}$$

rewritten as UPDATE FORMULA:

$$\boxed{X(t+dt) = X(t) + A(X(t), t) dt + D^{1/2}(X(t), t) N(t) \sqrt{dt}} \quad (4)$$

$D=0$: continuous deterministic Markov process (mem-loss)

$D \neq 0$: continuous stochastic Markov process (mem-loss)

X is continuous but NOT differentiable since:

$$\frac{X(t+dt) - X(t)}{dt} = A(X(t), t) + D^{1/2}(X(t), t) \frac{N(t)}{\sqrt{dt}}$$

diverges as $dt \rightarrow 0$.

$$\frac{N(t)}{\sqrt{dt}} = \sqrt{dt}^{-1} N(0,1) = N(0, \frac{1}{dt})$$

We define GAUSSIAN WHITE NOISE:

$$\boxed{W(t) \equiv \lim_{dt \rightarrow 0} N(0, \frac{1}{dt})} \quad (5)$$

rewriting Equation (3)/(4)

$$\boxed{\frac{dX(t)}{dt} = A(X(t), t) + W(t) D^{1/2}(X, t)}$$

White Noise form of the Langevin Eq.

Properties:

$$\begin{aligned} \langle W(t) \rangle &= 0 \\ \langle W(t) W(t') \rangle &= \delta(t-t') \\ \langle X(t') W(t) \rangle &= 0 \text{ for all } t' \leq t \end{aligned}$$

- Side remark (comes later in class): $P(x, t)$ is the probability density function of the random variable satisfies the FORWARD FOKKER PLANCK Eqs:

$$\partial_t P(x, t) = -\partial_x [A(x(t), t) P(x, t)] + \frac{1}{2} \partial_x^2 [D(x, t) P(x, t)]$$

- Note also that $X(t)$ does not have a proper derivative. But

$$\frac{dY(t)}{dt} \equiv X(t) \text{ has since } Y(t+dt) = Y(t) + X(t)dt$$

The Ornstein-Uhlenbeck process

$$A(x,t) = -\frac{1}{T}x \quad D(x,t) = c$$

Important for Brownian motion, Johnson Noise

Note: $\frac{dv}{dt} = \frac{1}{m}F + \frac{F(t)}{m}$

Update formula: $X(t+dt) = X(t) - \frac{1}{T}X(t)dt + \sqrt{c}N(t)\sqrt{dt}$
 or $\frac{dX}{dt} = -\frac{1}{T}X(t) + \sqrt{c} \cdot N(t)$ standard Langevin equation

initial condition $X(t_0) = x_0$

(1) Take average $\langle X(t+dt) \rangle = \langle X(t) \rangle - \frac{1}{T}\langle X(t) \rangle dt \quad \langle N(t) \rangle = 0$

$$\Leftrightarrow \frac{d\langle X(t) \rangle}{dt} = -\frac{1}{T}\langle X(t) \rangle$$

$$\langle X(t) \rangle = e^{-(t-t_0)/T} \cdot x_0 \quad \text{MEAN}$$

Next we square ("ITO CALCULUS")

$$X^2(t+dt) = X^2(t) - \frac{2}{T}X(t)^2 dt + 2\sqrt{c}N(t)X(t)\sqrt{dt} + cN^2(t)dt$$

average: $\langle X^2(t+dt) \rangle = \langle X^2(t) \rangle - \frac{2}{T}\langle X^2(t) \rangle dt + c dt$
 $\Rightarrow \frac{d\langle X^2(t) \rangle}{dt} = -\frac{2}{T}\langle X^2(t) \rangle + c$

$$\text{var}\{X(t)\} = \langle \Delta X(t)^2 \rangle = \frac{cT}{2} (1 - e^{-2(t-t_0)/T}) \quad \text{VARIANCE}$$

Since $X(t)$ is normal we conclude that for $X(t_0) = x_0$

$$X(t) = N\left(x_0 e^{-(t-t_0)/T}; \frac{cT}{2} (1 - e^{-2(t-t_0)/T})\right)$$

i.e. for $t \gg T \quad N(0, \frac{cT}{2}) \quad \text{long time limit}$

• STATIONARY STATISTICAL PREDICTION

Analogous $\text{cov}\{X_1(t_1), X_2(t_2)\} = \frac{cT}{2} e^{-(t_2-t_1)/T} (1 - e^{-2(t_1-t_0)/T})$

We may also need sometimes the integral (remember Langevin Eqs is for the velocity not position)

with proof $\langle Y(t) \rangle = \int_{t_0}^t dt' \langle X(t') \rangle$; $\langle Y^2(t) \rangle = 2 \int_{t_0}^t \int_{t_0}^{t'} \langle X(t') X(t'') \rangle dt' dt''$

without proof

$$\langle Y(t) \rangle = x_0 (1 - e^{-(t-t_0)/T}) \rightarrow x_0 T$$

$$\langle Y^2 \rangle - \langle Y \rangle^2 = \text{VAR}\{Y(t)\} = c \cdot T^2 \left[\Delta t - 2(1 - e^{-\Delta t/T}) + \frac{1}{2}(1 - e^{-2\Delta t/T}) \right]$$

not a
Stable
Process

$\rightarrow cT^2 \Delta t$ NOT STATIONARY



Wiener Process $A(x,t)=0$ $DC(x,t)=c$

Driftless Wiener process. $X(t+dt) = X(t) + \sqrt{c} N(t) \sqrt{dt}$

$\Rightarrow X(t) = N(x_0, c(t-t_0))$ solution

i.e. the O-U process in the limit of $\tau \rightarrow \infty$ becomes a driftless Wiener process. Example: phase diffusion in a LASER

SPECTRAL DENSITY FUNCTION

One advantage of the Markovian description of noise is the (over thermodynamical ones) is the insights into the ^{freq.} spectrum of fluctuations.

$$G(\tau) = \langle X(t)X(t+\tau) \rangle = C_X(\tau) \quad \text{auto-covariance (or autocorrelation)}$$

Then:

$$\left[\begin{aligned} S_{XX}(\omega) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-i\omega\tau} C_X(\tau) d\tau \quad \text{spectrum } \frac{e^2 [m^2]}{Hz} \text{ with } C_X(\tau) \in \mathbb{R} \\ S_{XX}(\omega) &= \frac{1}{\pi} \int_0^{\infty} C_X(\tau) \cos(\omega\tau) d\tau \quad \text{spectral density} \end{aligned} \right]$$

ω : frequency
 ω : angular-frequency.

Conversely:

$$\left[\begin{aligned} G(\tau) = C_X(\tau) &= \int_{-\infty}^{\infty} e^{-i\omega\tau} S_{XX}(\omega) d\omega \\ S_{XX}(\omega) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-i\omega\tau} G(\tau) d\tau \end{aligned} \right]$$

WIENER KHINCHIN THEOREM

$\langle X^2(t) \rangle = \sigma^2$ INTENSITY

it follows

$$\langle X^2(t) \rangle = \int_0^{\infty} S_{XX}(\nu) d\nu$$

integral over the spectrum gives the VARIANCE

$S_{XX}(\nu) d\nu$ is portion of intensity with frequencies $\nu \dots \nu+d\nu$.

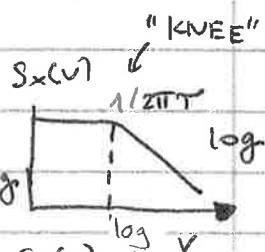
Example

O-U process:

$$C_X(t) = \frac{c\tau}{2} \cdot e^{-t/\tau}$$

$$S_{XX}(\nu) = \frac{2c\tau^2}{1+(2\pi\nu\tau)^2}$$

exponentially decaying
SPECTRAL DENSITY
spectrum of fluctuations

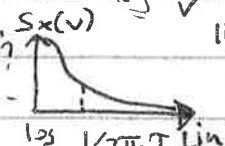


lin Second Example: Random Force / Function.

Gaussian white noise

$$\left[\begin{aligned} C_T(t) = \langle X_T(t)X_T(t) \rangle &= \sigma^2(t) \delta^2 \\ S_{TT}(\omega) &= \sigma^2 \int_{-\infty}^{\infty} \delta(t') e^{-i\omega t'} dt' = \sigma^2 \end{aligned} \right]$$

$\Rightarrow$ FLAT i.e. WHITE NOISE



Note spectral density of a Wiener process:

problem: $\langle x^2(t) \rangle_{\text{Wiener}} = \frac{\langle N(C(t-t_0))^2 \rangle}{x_0}$

$$= C(t-t_0) \quad \text{infinite intensity!}$$

But: $S_{xx}(\nu) = \frac{2CT^2}{1+(2\pi\nu T)^2} \quad (\text{for OU process})$

In the limit $T \rightarrow \infty$ (driftless Wiener Process):

$$S_{xx}(\nu) \rightarrow \frac{2C}{2\pi\nu^2} \propto \frac{1}{\nu^2} \quad (1/f^2 \text{ NOISE})$$

only valid for high frequencies since $(t-t_0) \gg T$
problem for $T \rightarrow \infty$

(L3)

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Finally back to Langevin:

$$F(t) = m \ddot{x}(t)$$

TIME EVOLUTION eq

$$\left[\frac{dy}{dt} = -\frac{1}{\tau} y(t) + \sqrt{c} \Gamma(t) \right]$$

$$\frac{1}{m} = \frac{1}{\tau}$$

- Thus, this is a O-U (Ornstein-Uhlenbeck) process with relaxation time τ and diffusion c . $\Rightarrow$ MEMORYLESS + ZERO MEAN
- Fixing c and τ : $\rightarrow$ s.e. + CONTINUOUS
- $\rightarrow$ need to be prop to Gaussian white noise

we have now:

$$\left[\begin{aligned} \langle V(t \rightarrow \infty) \rangle &= N(0, k_B T / m) \\ \langle x^2(t \rightarrow \infty) \rangle &= 2D(t - t_0) \end{aligned} \right]$$

Follow from Maxwell Boltzmann distribution
 $\sigma^2 = \Delta y^2 = k_B T / m$
 diffusion behavior from observation

$$\left[\begin{aligned} \langle V(t \rightarrow \infty) \rangle &= N(0, c\tau/2) \\ \langle x^2(t \rightarrow \infty) \rangle &= c\tau^2(t - t_0) \end{aligned} \right]$$

diffusion behavior from observation

from OU math

$$\left[D = \frac{k_B T}{m\gamma} \right]$$

$$\text{Thus } c\tau/2 = k_B T / m$$

$$\left[\tau = \frac{Dm}{k_B T} \right] \text{ or } \gamma = \frac{Dm}{k_B T}$$

Solution

$$V(t) = N\left(0, \frac{k_B T}{m} \left(1 - e^{-2(t-t_0)/\tau}\right)\right)$$

$$\text{OR } c = \frac{2}{D} \left(\frac{k_B T}{m}\right)^2$$

$$\text{Thus } \langle F(t) F(t') \rangle = \delta(t - t') 2k_B T \gamma$$

Hence, importantly we obtain

FLUCTUATION DISSIPATION THEOREM

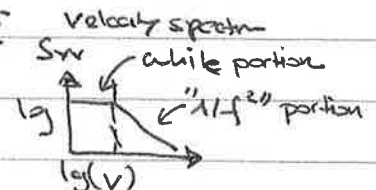
$$\gamma = \frac{1}{2k_B T} \int_{-\infty}^{\infty} \langle F(t) F(t+t') \rangle dt'$$

NOISE $\rightarrow$ DISSIPATION

$$\text{Also } \langle V(t) V(t+t') \rangle = \frac{k_B T}{m} e^{-|t'|/\tau}$$

Particle Fluctuation spectrum. $\langle V \cdot F \rangle = P = \frac{3}{2} k_B T$

$$P(\omega) = \frac{12 k_B T}{1 + \left(\frac{2\pi \omega m}{\gamma}\right)^2}; \quad S_{xx}(\omega) = \frac{k_B T / m \cdot 12}{1 + (2\pi m \omega \tau)^2}$$



(2) Johnson Noise 1928 by Johnson electrical noise + Nyquist

$$-R \cdot I(t) + V(t) = L \frac{dI(t)}{dt} = 0$$

key parameter R not γ

$$-L \frac{dI}{dt} = R \langle I(t) \rangle$$

$\langle V(t) \rangle = 0$
 Random voltage
 $L/R \equiv \tau$

* Fluct. voltage
 $[V = L \sqrt{c} \Gamma(t)]$

$$\Leftrightarrow \frac{dI}{dt} = -\frac{1}{\tau} I + \sqrt{c} \Gamma(t) \quad \text{or: } \frac{dy}{dt} = -\frac{1}{\tau} y + \sqrt{c} \gamma \Gamma(t)$$

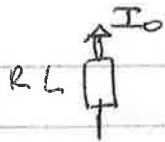
$$\text{Fixing } c \text{ by Boltzmann: } \left[k_B T / 2 = \left\langle \frac{1}{2} L I^2 \right\rangle \right]$$

$$\text{Thus: } c = \frac{2 k_B T R}{L}$$

$$\text{i.e. } V(t) = 2 k_B T R \Gamma(t)$$

Likewise

$$\left[R = \frac{1}{2k_B T} \int_{-\infty}^{\infty} V(t) V(t+t') dt' \right] \quad S_I = \frac{4k_B T}{R} \frac{1}{1 + \left(\frac{2\pi L \omega}{R}\right)^2}$$



$I(t_0) = i_0$: initial condition

$$I(t) = N \left(i_0 e^{-\frac{R}{L}(t-t_0)} ; \frac{K T}{L} (1 - e^{-\frac{2R}{L}(t-t_0)}) \right)$$